

MOdelling REvisited + MOdel REduction ERC-GZ project LL1202 - MORE





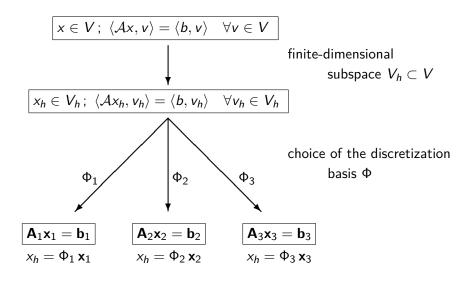
Interpretation of the algebraic error and algebraic preconditioning as the transformation of the discretization basis

Jan Papež, Josef Málek, Zdeněk Strakoš, Mattia Tani

Charles University in Prague Institute of Computer Science AS CR Universita di Bologna

September 23, 2014

Motivation



Motivation

Question 1

Can we interpret algebraic preconditioning as transformation of the discretization basis functions?

Question 2

Can we interpret algebraic error as transformation of the discretization basis functions such that the computed non-Galerkin solution is the Galerkin solution with respect to the transformed basis?

Outline

- 1 Setting and notation
- 2 Algebraic preconditioning and the transformation of the discretization basis
- 3 Interpretation of the algebraic error

1. Basic notation

Let $\,V\,$ be a real (infinite dimensional) Hilbert space with the inner product

$$(\cdot,\cdot)_V:V\times V\to \mathbf{R}$$
,

 $V^{\#}$ be the dual space of bounded linear functionals on V with the duality pairing

$$\langle \cdot, \cdot \rangle : V^{\#} \times V \to \mathbf{R}$$
.

For each $f \in V^{\#}$ there exists a unique $\tau f \in V$ such that

$$\langle f, v \rangle = (\tau f, v)_V$$
 for all $v \in V$.

In this way the inner product $(\cdot,\cdot)_V$ determines the Riesz map

$$\tau: V^{\#} \to V$$
.

1. Functional formulation

Consider a PDE problem described in the form of the functional equation

$$Ax = b$$
, $A: V \to V^{\#}$, $x \in V$, $b \in V^{\#}$ (1)

where the linear, bounded, and coercive operator $\mathcal A$ is self-adjoint w.r.t. the duality pairing $\langle\cdot,\cdot\rangle$.

With the transformation using the Riesz map,

$$\tau Ax = \tau b$$
, $\tau A : V \to V$, $x \in V$, $\tau b \in V$,

which is called operator preconditioning.

Algebraic preconditioning and the transformation of the discretization basis

2. Related ideas and work

Operator preconditioning

```
[Klawonn (1995, 1996)]; [Arnold, Falk, and Winther (1997, 1997)]; [Steinbach and Wendland (1998)]; [McLean and Tran (1997)]; [Christiansen and Nédélec (2000, 2000)]; [Powell and Silvester (2003)]; [Elman, Silvester, and Wathen (2005, 2014)]; [Hiptmair (2006)]; [Axelsson and Karátson (2009)]; [Mardal and Winther (2011)]; [Kirby (2011)]; [Zulehner (2011)]; . . .
```

CG in Hilbert spaces

```
[Hayes (1954)]; ...; [Glowinski (2003)]; [Axelsson and Karátson (2009)]; [Mardal and Winther (2011)]; [Günnel, Herzog and Sachs (2013)].
```

Hierarchical preconditioning, multilevel methods, domain decomposition with coarse components

```
[Yserentant (1985, 1986)], using the work of [Zienkiewicz et al. (1970)]; ... [Axelsson and Vassilevski (1989)], the survey in [Vassilevski (2008)]; ... [Jarošová, Klawonn and Rheinbach (2012)]; ...
```

2. Krylov manifolds in Hilbert spaces

Using the Riesz map, one can form for $\ x_0 \in V$, $r_0 \equiv b - \mathcal{A} x_0 \in V^\#$ the Krylov subspace

$$K_n \equiv \operatorname{span}\{\tau r_0, (\tau \mathcal{A})(\tau r_0), (\tau \mathcal{A})^2(\tau r_0), \dots, (\tau \mathcal{A})^{n-1}(\tau r_0)\},\,$$

and approximate the solution $x = (\tau A)^{-1} \tau b$ constructing the approximations $x_n \in x_0 + K_n$ from the Krylov manifolds in V.

Looking for the approximate solution minimizing the energy leads to

$$||x - x_n||_a = \min_{z \in x_0 + K_n} ||x - z||_a, \qquad ||z||_a^2 \equiv \langle Az, z \rangle,$$

which is equivalent to the Galerkin orthogonality condition

$$\langle b - Ax_n, w \rangle = 0$$
 for all $w \in K_n$.

2. Preconditioned CG in Hilbert spaces

$$r_{0} = b - Ax_{0} \in V^{\#}, \quad p_{0} = \tau r_{0} \in V$$

$$x_{n} = x_{n-1} + \alpha_{n-1}p_{n-1},$$

$$\alpha_{n-1} = \frac{\langle r_{n-1}, \tau r_{n-1} \rangle}{\langle Ap_{n-1}, p_{n-1} \rangle} = \frac{(\tau r_{n-1}, \tau r_{n-1})_{V}}{(\tau Ap_{n-1}, p_{n-1})_{V}},$$

$$r_{n} = r_{n-1} - \alpha_{n-1}Ap_{n-1},$$

$$p_{n} = \tau r_{n} + \beta_{n}p_{n-1},$$

$$\beta_{n} = \frac{\langle r_{n}, \tau r_{n} \rangle}{\langle r_{n-1}, \tau r_{n-1} \rangle} = \frac{(\tau r_{n}, \tau r_{n})_{V}}{(\tau r_{n-1}, \tau r_{n-1})_{V}}.$$

2. Finite dimensional CG

$$\begin{split} \Phi &= \{\phi_1, \dots, \phi_N\} \\ &\quad \text{basis of the finite-dimensional subspace} \quad V_h \subset V \,, \\ \Phi^\# &= \{\phi_1^\#, \dots, \phi_N^\#\} \\ &\quad \text{canonical basis of the dual} \quad V_h^\# \,, \quad \Phi^\# \Phi = \mathbf{I} \,. \\ \text{Using the coordinates in} \quad \Phi \quad \text{and in} \quad \Phi^\# \,, \\ &\quad \langle f, v \rangle = \quad \langle \Phi^\# \mathbf{f}, \Phi \mathbf{v} \rangle \quad = \mathbf{v}^* \mathbf{f} \,, \\ &\quad \langle f, v \rangle = \quad \langle \Phi^\# \mathbf{f}, \Phi \mathbf{v} \rangle \quad = \mathbf{v}^* \mathbf{M} \mathbf{u} \,, \\ &\quad \langle f, v \rangle_{V_h} = \quad \langle \Phi \mathbf{u}, \Phi \mathbf{v} \rangle_{V_h} \quad = \mathbf{v}^* \mathbf{M} \mathbf{u} \,, \\ &\quad \mathcal{A} u = \quad \mathcal{A} \Phi \, \mathbf{u} \quad = \Phi^\# \mathbf{A} \mathbf{u} \,, \\ &\quad \mathcal{T} f = \quad \mathcal{T} \Phi^\# \mathbf{f} \quad = \Phi \, \mathbf{M}^{-1} \, \mathbf{f} \,; \end{split}$$
 where
$$\mathbf{M} = [M_{ij}] \quad = \quad [(\phi_j, \phi_i)_{V_h}] \,, \end{split}$$

 $\mathbf{A} = [A_{ii}] = [\langle \mathcal{A}\phi_i, \phi_i \rangle], \quad i, j = 1, \dots, N.$

2. Preconditioned algebraic CG

Preconditioned CG

With $b = \Phi^{\#} \mathbf{b}$, $x_n = \Phi \mathbf{x}_n$, $p_n = \Phi \mathbf{p}_n$, $r_n = \Phi^{\#} \mathbf{r}_n$ we get the standard preconditioned algebraic CG with the preconditioner \mathbf{M} .

Unpreconditioned CG is in this setting an oxymoron!

Unpreconditioned CG, i.e. $\mathbf{M}=\mathbf{I}$, corresponds to the basis Φ orthonormal w.r.t. the inner product $(\cdot,\cdot)_{V_h}$.

2. Orthogonalization of the discretization basis

Consider the decomposition $\mathbf{M} = \mathbf{L}\mathbf{L}^*$ and the transformed discretization basis $\Phi_t = \Phi\left(\mathbf{L}^*\right)^{-1}$ that is orthonormal w.r.t. $(\cdot, \cdot)_{V_h}$.

Then the matrix representation of CG in the Hilbert space V_h with the basis Φ_t gives the standard algebraic (unpreconditioned) CG applied to the preconditioned system

$$L^{-1}A(L^*)^{-1}(L^*x) = L^{-1}b$$
.

2. Interpretation of the algebraic preconditioning

For the algebraic preconditioning with $\widehat{\mathbf{L}}\widehat{\mathbf{L}}^* = \widehat{\mathbf{M}} \neq \mathbf{M}$, the (transformed) discretization basis $\widehat{\Phi} = \Phi(\widehat{\mathbf{L}}^*)^{-1}$ is not orthonormal w.r.t. $(\cdot,\cdot)_{V_h}$.

In order to obtain the interpretation of the algebraic preconditioning $\widehat{\mathbf{M}}$ as the transformation of the basis $\Phi \to \widehat{\Phi}$, we have to change also the inner product in V_h :

$$(u, v)_{V_h} = (\Phi \mathbf{u}, \Phi \mathbf{v})_{V_h} = \mathbf{v}^* \mathbf{M} \mathbf{u},$$

has to be replaced by

$$(u,v)_{\mathsf{new},V_h} = (\widehat{\boldsymbol{\Phi}}\,\widehat{\mathbf{u}},\widehat{\boldsymbol{\Phi}}\,\widehat{\mathbf{v}})_{\mathsf{new},V_h} \equiv \widehat{\mathbf{v}}^*\widehat{\mathbf{u}} = \mathbf{v}^*\widehat{\mathbf{M}}\mathbf{u}\,.$$

2. Observations

Algebraic preconditioning associated with the operator preconditioning is equivalent to the orthogonalization of the discretization basis in the given finite-dimensional Hilbert space V_h .

Algebraic preconditioning can be interpreted as transformation of the discretization basis and, at the same time, transformation of the inner product in V_h such that the transformed basis $\widehat{\Phi}$ is orthonormal with respect to the transformed inner product.

Interpretation of the algebraic error

3. Algebraic backward error

Let x solves the discretized algebraic system

$$\mathbf{A}\mathbf{x} = \mathbf{b}$$
, $A_{ij} = \langle \mathcal{A}\phi_j, \phi_i \rangle$, $b_i = \langle b, \phi_i \rangle$

and let the algebraic vector $\hat{\mathbf{x}}$ which approximates \mathbf{x} solves the perturbed system

$$(\mathbf{A} + \mathbf{E}) \, \widehat{\mathbf{x}} = \mathbf{b} + \mathbf{f} \, .$$

Our aim is to interpret the perturbations \mathbf{E}, \mathbf{f} as transformation of the discretization bases.

We will consider possibly different transformation of the discretization (search) and the test bases

$$\begin{split} \widehat{\boldsymbol{\Phi}} &= & \boldsymbol{\Phi} \left(\boldsymbol{I} + \boldsymbol{D} \right), \\ \widetilde{\boldsymbol{\Phi}} &= & \boldsymbol{\Phi} \left(\boldsymbol{I} + \boldsymbol{G} \right). \end{split}$$

$$\widetilde{\Phi} = \Phi(\mathbf{I} + \mathbf{G})$$

3. Transformation of the discretization bases

Let the Galerkin solution $x_h = \Phi x$ of the discretized formulation

$$\langle \mathcal{A}x_h, \phi_i \rangle = \langle b, \phi_i \rangle, \quad i = 1, \dots, N$$

can be expressed as the Galerkin solution $x_h = \hat{\Phi} \hat{\mathbf{x}} = \Phi (\mathbf{I} + \mathbf{D}) \hat{\mathbf{x}}$ of the formulation

$$\langle \mathcal{A}x_h, \widetilde{\phi}_i \rangle = \langle b, \widetilde{\phi}_i \rangle, \quad i = 1, \dots, N.$$

Discretization via the transformed basis functions $\widehat{\Phi}$ and the test functions $\widetilde{\Phi}$ results in the linear algebraic system

$$\overline{\mathbf{A}}\widehat{\mathbf{x}} = \overline{\mathbf{b}}, \qquad \overline{A}_{ij} = \langle \mathcal{A}\widehat{\phi}_j, \widetilde{\phi}_i \rangle, \quad \overline{b}_i = \langle b, \widetilde{\phi}_i \rangle,$$

 $\text{where} \quad \overline{\mathbf{A}} = (\mathbf{I} + \mathbf{G})^T \mathbf{A} \, (\mathbf{I} + \mathbf{D}) \quad \text{and} \quad \overline{\mathbf{b}} = (\mathbf{I} + \mathbf{G})^T \mathbf{b} \, .$

3. Interpretation of the algebraic error

The algebraic vector $\hat{\mathbf{x}}$ which approximates \mathbf{x} solves exactly the algebraic system determined by the Galerkin discretization of the infinite dimensional problem.

Interpreting the perturbations in

$$(\mathbf{A} + \mathbf{E})\,\widehat{\mathbf{x}} = \mathbf{b} + \mathbf{f}$$

as transformation of the discretization bases gives

$$\mathbf{A} + \mathbf{E} = (\mathbf{I} + \mathbf{G})^T \mathbf{A} (\mathbf{I} + \mathbf{D}),$$

$$\mathbf{b} + \mathbf{f} = (\mathbf{I} + \mathbf{G})^T \mathbf{b}.$$

3. Three classes of backward errors

Perturbed algebraic system

$$\left(\boldsymbol{A}+\boldsymbol{E}\right)\widehat{\boldsymbol{x}}=\boldsymbol{b}+\boldsymbol{f}\,.$$

- **1 E** \neq 0, **f** \neq 0,
- **2** $\mathbf{A} + \mathbf{E}$ is symmetric positive definite,
- **3** f = 0.

3. General case ($\mathbf{E} \neq 0$, $\mathbf{f} \neq 0$)

Given perturbations ${\bf E},{\bf f}$ satisfying $({\bf A}+{\bf E}){\bf \widehat x}={\bf b}+{\bf f}$, there are infinitely many couples of matrices ${\bf D},{\bf G}$ satisfying

$$\begin{aligned} \mathbf{A} + \mathbf{E} &= & \left(\mathbf{I} + \mathbf{G} \right)^T \mathbf{A} \left(\mathbf{I} + \mathbf{D} \right), \\ \mathbf{b} + \mathbf{f} &= & \left(\mathbf{I} + \mathbf{G} \right)^T \mathbf{b}. \end{aligned}$$

Pick any \mathbf{G} such that $\mathbf{G}^T\mathbf{b} = \mathbf{f}$ and $(\mathbf{I} + \mathbf{G})$ is nonsingular. Set $\mathbf{D} = \mathbf{A}^{-1} \left(\mathbf{I} + \mathbf{G}^T \right)^{-1} \left(\mathbf{E} - \mathbf{G}^T \mathbf{A} \right) \,.$

Outlook: The optimal choice of D, G?

3. Symmetric case (A + E is SPD)

In order to preserve the symmetry take $\, {f D} = {f G} \, .$ Then

$$\mathbf{A} + \mathbf{E} = (\mathbf{I} + \mathbf{D})^{T} \mathbf{A} (\mathbf{I} + \mathbf{D}), \qquad (2)$$

$$\mathbf{b} + \mathbf{f} = (\mathbf{I} + \mathbf{D})^T \mathbf{b}. \tag{3}$$

D satisfies (2) if and only if

$$\mathbf{I} + \mathbf{D} = \mathbf{A}^{-1/2} \mathbf{U} (\mathbf{A} + \mathbf{E})^{1/2}$$
, \mathbf{U} is an orthogonal matrix.

From (3) we obtain the condition

$$\mathbf{b}^T\mathbf{x} = (\mathbf{b} + \mathbf{f})^T\widehat{\mathbf{x}}\,, \quad \text{or, equivalently,} \quad \|\mathbf{x}\|_{\mathbf{A}} = \|\widehat{\mathbf{x}}\|_{\mathbf{A} + \mathbf{E}}\,,$$

which in general does not hold.

3. Backward error with f = 0

When $\mathbf{f}=0$, it is natural to set $\mathbf{G}=0$, i.e. to consider original test functions. This case was considered in [Gratton, Jiránek, and Vasseur (2013)]; [P., Liesen, Strakoš (2014)].

From
$$\mathbf{A} + \mathbf{E} = \mathbf{A} \left(\mathbf{I} + \mathbf{D} \right)$$
 we have $\mathbf{A} \mathbf{D} = \mathbf{E}$ and $\mathbf{D} = \mathbf{A}^{-1} \mathbf{E}$.

Loss of locality: the transformed basis $\widehat{\Phi} = \Phi(\mathbf{I} + \mathbf{D})$ has global support (\mathbf{D} is dense)!

3. Loss of locality

Illustration on 1D Poisson model problem, using two perturbations in

$$(\mathbf{A} + \mathbf{E}) \, \widehat{\mathbf{x}} = \mathbf{b} \,.$$

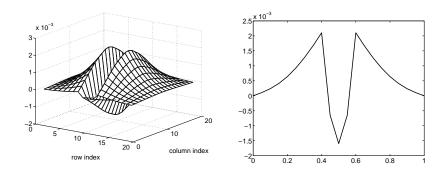
•
$$\mathbf{E} = \frac{(\mathbf{b} - \mathbf{A}\widehat{\mathbf{x}})\widehat{\mathbf{x}}^T}{\|\widehat{\mathbf{x}}\|_2^2}$$
, then $\mathbf{D} = \mathbf{A}^{-1}\mathbf{E} = \frac{(\mathbf{x} - \widehat{\mathbf{x}})\widehat{\mathbf{x}}^T}{\|\widehat{\mathbf{x}}\|_2^2}$,

ullet symmetric perturbation $oldsymbol{\mathsf{E}}_{\mathsf{sym}}$ with the minimal Frobenius norm

$$\mathbf{E}_{\mathrm{sym}} = \arg\min\left\{ \|\mathbf{E}\|_F \mid \mathbf{E} = \mathbf{E}^T, \ (\mathbf{A} + \mathbf{E}) \, \widehat{\mathbf{x}} = \mathbf{b}
ight\};$$

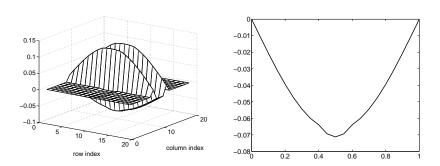
see [Bunch, Demmel, and van Loan (1989)].

3. Transformed basis



MATLAB surf plot of the transformation matrix $\mathbf{D} = \mathbf{A}^{-1}\mathbf{E}$ (left) and the difference $\hat{\phi}_j - \phi_j$ (right).

3. Transformed basis for symmetric perturbation matrix



MATLAB surf plot of the transformation matrix $\mathbf{D} = \mathbf{A}^{-1}\mathbf{E}_{\text{sym}}$ (left) and the difference $\hat{\phi}_j - \phi_j$ (right).

Please note that $\|\mathbf{A}^{-1}\mathbf{E}_{svm}\| \gg \|\mathbf{A}^{-1}\mathbf{E}\|!$

References



J. Málek and 7. Strakoš:

Preconditioning and the Conjugate Gradient Method in the Context of Solving PDEs.

SIAM Spotlight Series, SIAM, Philadelphia, (2014). (in print)



Energy backward error: interpretation in numerical solution of elliptic partial differential equations and behaviour in the conjugate gradient method.

Electron. Trans. Numer. Anal. 40 (2013), pp. 338–355.



J. Papež, J. Liesen, Z. Strakoš:

Distribution of the discretization and algebraic error in numerical solution of partial differential equations.

Linear Algebra Appl. 449 (2014), pp. 89–114.

Thank you for your attention!