# <span id="page-0-0"></span>Measure-valued solutions to compressible Euler system

### Piotr Gwiazda

## joint work with Eduard Feireisl, Ondřej Kreml, Agnieszka Swierczewska-Gwiazda and Emil Wiedemann ´

University of Warsaw and Institute of Mathematics Polish Academy of Sciences

Roztoky, 3rd April 2017

# Plan of the talk

- Measure-valued strong uniqueness to compressible Euler and Navier-Stokes like systems
	- P. G, A. Swierczewska.-Gwiazda., E. Wiedemann, ´ Weak-Strong Uniqueness for Measure-Valued Solutions of Some Compressible Fluid Models, Nonlinearity, 2015
	- E. Feireisl, P. G, A. Swierczewska.-Gwiazda., E. Wiedemann, ´ Dissipative measure-valued solutions to the compressible Navier-Stokes system, Calculus of Variations and Partial Differential Equations, 2016
- Measure-valued strong uniqueness for general hyperbolic conservation laws
	- . P. G, O. Kreml, A. Świerczewska.-Gwiazda, Dissipative measure-valued solutions for general hyperbolic conservation laws, preprint

This system models the flow of an inviscid, incompressible fluid with constant density in the absence of external forces. The vector-valued function  $v(t, x)$  is the velocity of the fluid and the scalar-valued function  $p(t, x)$  is the pressure

$$
v_t + \operatorname{div} (v \otimes v) + \nabla p = 0,
$$
  
div  $v = 0$ .

Leonhard Euler, 1757

For isentropic Euler equations, in contrast to the incompressible system, the pressure  $p = p(\rho)$  is no longer a Lagrange multiplier, but a constitutively given function of the density

$$
\partial_t(\rho v) + \operatorname{div}(\rho v \otimes v) + \nabla p(\rho) = 0, \n\partial_t \rho + \operatorname{div}(\rho v) = 0.
$$

Consider the equations of nonlinear elasticity

$$
\partial_t F = \nabla_x \mathbf{u}
$$
  

$$
\partial_t \mathbf{u} = \text{div}_x (D_F W(F))
$$

for an unknown matrix field  $F$  and an unknown vector field  $\mathbf{u}$ . Function  $W(F)$  is given. The associated entropy is given by

$$
\eta(\mathbf{u},F)=\frac{1}{2}|\mathbf{u}|^2+W(F).
$$

# Notation

 $C_0(\mathbb{R}^d)$  – closure of the space of continuous functions on  $\mathbb{R}^d$  with compact support w.r.t. the  $\|\cdot\|_{\infty}$ -norm.  $(\mathcal{C}_0(\mathbb{R}^d))^* \cong \mathcal{M}(\mathbb{R}^d)$  — the space of signed Radon measures with finite mass. The duality pairing is given by

$$
\langle \mu, f \rangle = \int\limits_{\mathbb{R}^d} f(\lambda) \, d\mu(\lambda).
$$

#### **Definition**

A map  $\mu:\Omega\to\mathcal{M}(\mathbb{R}^d)$  is called weakly\* measurable if the functions  $x \mapsto \langle \mu(x), f \rangle$  are measurable for all  $f \in \mathcal{C}_0(\mathbb{R}^d).$ 

# Young measures

- The main feature of Young measure theory is that it allows us to pass to a limit in the expression  $f(\mathbf{\nu}^j)$  with nonlinear  $f$  and only weakly-star convergent  $v^j$ .
- What's the strategy? Instead of considering  $f(v^j)$  we embed the problem in a larger space, but gain linearity, i.e.  $\langle f, \delta_{\mathsf{v}^j(\mathsf{x})}\rangle.$
- If  $f \in C_0(\mathbb{R})$  using the duality

$$
(L^1(\Omega; C_0(\mathbb{R}^d)))^* \cong L^\infty_w(\Omega; \mathcal{M}(\mathbb{R}^d)),
$$

Banach-Alaoglu theorem and weak-star continuity of linear operators allows for limit passage to get  $\langle f, \nu_x \rangle$ .

• What's the cost?

We end up with a weaker notion of solution: *measure-valued* solution

# Can the Young measures describe a concentration effect?

#### **Definition**

A bounded sequence  $\{z^j\}$  in  $L^1(\Omega)$  converges in biting sense to a function  $z\in L^1(\Omega)$ , written  $z^j\stackrel{b}{\to}z$  in  $\Omega$ , provided there exists a sequence  ${E_k}$  of measurable subsets of  $\Omega$ , satisfying  $\lim_{k\to\infty} |E_k| = 0$ , such that for each k

$$
z^j \rightharpoonup z \quad \text{in} \quad L^1(\Omega \setminus E_k).
$$

#### Remarks

Biting limit can be also express as  $\lim_{n\to\infty}\lim_{j\to\infty} T^n(z^j)$ , where by  $T^{n}(\cdot)$  we denote standard truncation operator.

#### Lemma

Let  $u^j$  be a sequence of measurable functions and  $\nu_x$  a Young measure associated to a subsequence  $u_{j_k}$ . Then  $f(\cdot, u^{j_k}) \stackrel{b}{\to} \langle \nu_x, f \rangle$ for every Carathéodory function  $f(\cdot, \cdot)$  s.t.  $f(\cdot, \mu^{j_k})$  is a bounded sequence in  $L^1(\Omega)$ . Here  $\langle \nu_{\mathsf{x}}, f \rangle = \int_{\mathbb{R}^d} f \, d\nu_{\mathsf{x}}$ .

#### Remarks

In view of the above facts the classical Young measures prescribe only the oscillation effect, not the concentration one. The attempt to prescribe also concentration effect by some generalizations of the Young measures was initiated by DiPerna and Majda

#### Lemma

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• Y. Brenier, C. De Lellis, L. Székelyhidi, Jr., Weak-strong uniqueness for measure-valued solutions. Comm. Math. Phys. 2011,

Incompressible Euler -oscillation and concentration measure,

S. Demoulini, D. M. A. Stuart, A. E. Tzavaras, Weak-strong uniqueness of dissipative measure-valued solutions for polyconvex elastodynamics. Arch. Ration. Mech. Anal. 2012 In weak formulation only oscillation measure, in entropy inequality there appears non-negative concentration measure

- P. G, A. Swierczewska-Gwiazda, E. Wiedemann, Weak-Strong ´ Uniqueness for Measure-Valued Solutions of Some Compressible Fluid Models, Nonlinearity, 2015 Oscillatory and vector-valued concentration measure both in weak formulation and entropy inequality
- E. Feireisl, P. G, A. Swierczewska-Gwiazda, E. Wiedemann, ´ Dissipative measure-valued solutions to the compressible Navier–Stokes system, Calculus of Variations and Partial Differential Equations, 2016

Instead of vector-valued concentration measure the dissipation defect is introduced

J. Brezina, E. Feireisl, Measure-valued solutions to the complete Euler system, arXiv:1702.04870

- $\nu_{\mathsf{x},t} \in \mathcal{P}(\mathbb{R}^d)$  for a.e.  $(\mathsf{x},t) \in \mathbb{R}^d \times \mathbb{R}^+$  (oscillation measure)
- $m \in \mathcal{M}^{+}(\mathbb{R}^{d}\times\mathbb{R}^{+})$  (concentration measure)
- $\nu_{\mathsf{x},t}^\infty \in \mathcal{P}(\mathcal{S}^{d-1})$  for *m*-a.e.  $(\mathsf{x},t) \in \mathbb{R}^d \times \mathbb{R}^+$ (concentration-angle measure)

R. J. DiPerna and A. J. Majda, Oscillations and concentrations in weak solutions of the incompressible fluid equations, Comm. Math. Phys. 1987.

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# Measure-valued solutions to incompressible Euler system

We say that  $(\nu, m, \nu^{\infty})$  is a measure-valued solution of IE with initial data  $\mathit{u}_0$  if for every  $\phi \in \mathcal{C}^1_{c,\mathrm{div}}([0,\,\mathcal{T})\times \mathbb{T}^n;\mathbb{R}^n)$  it holds that

$$
\int_0^T \int_{\mathbb{T}^n} \partial_t \phi \cdot \overline{u} + \nabla \phi : \overline{u \otimes u} dx dt + \int_{\mathbb{T}^n} \phi(\cdot, 0) \cdot u_0 dx = 0.
$$

**Where** 

$$
\overline{u} = \langle \lambda, \nu \rangle
$$
  

$$
\overline{u \otimes u} = \langle \lambda \otimes \lambda, \nu \rangle + \langle \beta \otimes \beta, \nu^{\infty} \rangle m
$$

If the solution is generated by some approximation sequences, then the black terms on right-hand side correspond to the biting limit of sequences whereas the blue ones corespond to concentration measure

Let us set

$$
E_{\mathsf{mvs}}(t):=\int_{\mathbb{T}^n}\frac{1}{2}\overline{|u|^2}(t,x)dx
$$

for almost every  $t$ , where

$$
\overline{|u|^2} = \langle |\lambda|^2, \nu \rangle + \langle |\beta|^2, \nu^{\infty} \rangle m
$$

and

$$
E_0 := \int_{\mathbb{T}^n} \frac{1}{2} |u_0|^2(x) dx.
$$

We then say that a measure-valued solution is admissible if

$$
E_{\text{mvs}}(t) \leq E_0
$$

in the sense of distributions.

#### Theorem (Y. Brenier, C. De Lellis, L. Székelyhidi, Jr., 2011)

Let  $U \in C^1([0,T] \times \mathbb{T}^n)$  be a solution of IE . If  $(\nu,m,\nu^\infty)$  is an admissible measure-valued solution with the same initial data, then

$$
\nu_{t,x} = \delta_{U(t,x)} \text{ for a.e. } t, x, \text{ and } m = 0.
$$

#### Remark:

Some generalization of this result: Emil Wiedemann, Weak-strong uniqueness in fluid dynamics, arXiv:1705.04220

Let's define relative energy (entropy):

$$
E_{rel}(t):=\int_{\mathbb{T}^n}\frac{1}{2}\overline{|u-U|^2}(t,x)dx
$$

where

$$
\overline{|u-U|^2}=\langle |\lambda-U|^2,\nu\rangle+\langle |\beta|^2,\nu^\infty\rangle m
$$

then

$$
\frac{d}{dt}E_{rel}(t)\leq ||U||_{C^1}\cdot E_{rel}(t).
$$

We say that  $(\nu, m, \nu^{\infty})$  is a measure-valued solution of CE with initial data  $(\rho_0, u_0)$  if for every  $\tau \in [0,\, T],\, \psi \in C^1([0,\, T]\times {\mathbb T}^n; {\mathbb R}),$  $\phi \in \mathcal{C}^1([0,\, \mathcal{T}] \times \mathbb{T}^n; \mathbb{R}^n)$  it holds that

$$
\int_0^{\tau} \int_{\mathbb{T}^n} \partial_t \psi \overline{\rho} + \nabla \psi \cdot \overline{\rho u} dxdt + \int_{\mathbb{T}^n} \psi(x,0)\rho_0 - \psi(x,\tau)\overline{\rho}(x,\tau)dx = 0,
$$

$$
\int_0^{\tau} \int_{\mathbb{T}^n} \partial_t \phi \cdot \overline{\rho u} + \nabla \phi : \overline{\rho u \otimes u} + \text{div } \phi \overline{\rho^{\gamma}} dxdt + \int_{\mathbb{T}^n} \phi(x,0) \cdot \rho_0 u_0 - \phi(x,\tau) \cdot \overline{\rho u}(x,\tau)dx = 0.
$$

#### **Where**

$$
\overline{\rho} = \langle \lambda_1, \nu \rangle
$$
  
\n
$$
\overline{\rho^{\gamma}} = \langle \lambda_1^{\gamma}, \nu \rangle + \langle \beta_1^{\gamma}, \nu^{\infty} \rangle m
$$
  
\n
$$
\overline{\rho u} = \langle \sqrt{\lambda_1} \lambda', \nu \rangle
$$
  
\n
$$
\overline{\rho u \otimes u} = \langle \lambda' \otimes \lambda', \nu \rangle + \langle \beta' \otimes \beta', \nu^{\infty} \rangle m
$$
  
\n
$$
\overline{\rho |u|^2} = \langle |\lambda'|^2, \nu \rangle + \langle |\beta'|^2, \nu^{\infty} \rangle m
$$

If the solution is generated by some approximation sequences, then the black terms on right-hand side correspond to the biting limit of sequences whereas the **blue** ones corespond to concentration measure

#### Note that:

$$
\begin{aligned} \|\langle \beta' \otimes \beta', \nu^{\infty} \rangle m\|_{\text{TV}} &\leq C \|\langle \text{tr}(\beta' \otimes \beta'), \nu^{\infty} \rangle m\|_{\text{TV}} \\ &= C \|\langle |\beta'|^2, \nu^{\infty} \rangle m\|_{\text{TV}} \end{aligned}
$$

Let us set

$$
E_{\mathsf{mvs}}(t):=\int_{\mathbb{T}^n}\frac{1}{2}\overline{\rho|u|^2}(t,x)+\frac{1}{\gamma-1}\overline{\rho^{\gamma}}(t,x)dx
$$

for almost every t, and

$$
E_0:=\int_{\mathbb{T}^n}\frac{1}{2}\rho_0|u_0|^2(x)+\frac{1}{\gamma-1}\rho_0^\gamma(x)dx.
$$

We then say that a measure-valued solution is admissible if

$$
E_{\text{mvs}}(t) \leq E_0
$$

in the sense of distributions.

## Theorem (G., Swierczewska-Gwiazda, Wiedemann, 2015) ´

Let  $R\in W^{1,\infty}([0,\,T]\times\mathbb{T}^n),$   $U\in C^1([0,\,T]\times\mathbb{T}^n)$  is a solution of CE with initial data  $\varrho_0 \geq c > 0$ ,  $\varrho_0 \in L^{\gamma}(\mathbb{T}^n)$ ,  $\varrho_0 |u_0|^2 \in L^1(\mathbb{T}^n)$ , and  $R(x, t) \ge c > 0$  for some constant c and all  $(t, x) \in [0, T] \times \mathbb{T}^n$ . If  $(\nu, m, \nu^{\infty})$  is an admissible measure-valued solution with the same initial data, then

$$
\nu_{t,x} = \delta_{(R(t,x),\sqrt{R(t,x)})} U(t,x)
$$
 for a.e. t, x, and m = 0.

#### Remark

An analogue result holds for slightly weaker notion of solutions

## General hyperbolic conservation law

• Y. Brenier, C. De Lellis, L. Székelyhidi, Jr., Weak-strong uniqueness for measure-valued solutions. Comm. Math. Phys. 2011,

General hyperbolic systems - only oscillation measure, both in weak formulation and entropy inequality

S. Demoulini, D. M. A. Stuart, A. E. Tzavaras, Weak-strong uniqueness of dissipative measure-valued solutions for polyconvex elastodynamics. Arch. Ration. Mech. Anal. 2012 In weak formulation only oscillation measure, in entropy inequality there appears non-negative concentration measure

## General hyperbolic conservation law

- C. Christoforou, A. Tzavaras, Relative entropy for hyperbolic-parabolic systems and application to the constitutive theory of thermoviscoelasticity, arXiv:1603.08176 An analogue result for more general form of a system, hyperbolic-parabolic case, also only with a non-negative concentration measure in entropy inequality
- P. G., O. Kreml, A. Swierczewska-Gwiazda. Dissipative ´ measure valued solutions for general hyperbolic conservation laws Concentration measure both in the weak formulation and the entropy inequality

We consider the hyperbolic system of conservation laws in the form

<span id="page-27-0"></span>
$$
\partial_t A(u) + \partial_\alpha F_\alpha(u) = 0 \tag{1}
$$

with the initial condition  $\mathit{u}(0) = \mathit{u}_0.$  Here  $\mathit{u}:[0,\mathcal{T}]\times\mathbb{T}^d\rightarrow\mathbb{R}^n.$ There exists an open convex set  $X \subset \mathbb{R}^n$  such that the mappings  $A:\overline{X}\to \mathbb{R}^n$ ,  $\mathsf{F}_\alpha:X\to \mathbb{R}^n$  are  $\mathcal{C}^2$  maps on  $X$ ,  $A$  is continuous on  $\overline{X}$  and  $\nabla A(u)$  is nonsingular for all  $u \in X$ .

# Hypotheses

(H1) There exists an entropy-entropy flux pair  $(\eta, q_\alpha)$ ,  $\eta(u) \geq 0$ and lim  $\eta(u) = \infty$  $|u| \rightarrow \infty$ This yields the existence of a  $\mathcal{C}^1$  function  $\mathcal{G}:\overline{X}\to\mathbb{R}^n$  such that

$$
\nabla \eta = G \cdot \nabla A, \quad \nabla q_{\alpha} = G \cdot \nabla F_{\alpha}, \qquad \alpha = 1, ..., d.
$$

(H2) The symmetric matrix

$$
\nabla^2 \eta(u) - G(u) \cdot \nabla^2 A(u)
$$

is positive definite for all  $u \in X$ .

(H3) The vector  $A(u)$  and the fluxes  $F_\alpha(u)$  are bounded by the entropy, i.e.

$$
|A(u)|\leq C(\eta(u)+1),\quad |F_{\alpha}(u)|\leq C(\eta(u)+1),\qquad \alpha=1,...,d.
$$

Define for a strong solution U taking values in a compact set  $D \subset X$  the relative entropy

$$
\eta(u|U) := \eta(u) - \eta(U) - \nabla \eta(U) \cdot \nabla A(U)^{-1}(A(u) - A(U))
$$
  
= 
$$
\eta(u) - \eta(U) - G(U) \cdot (A(u) - A(U))
$$

and the relative flux as

$$
F_{\alpha}(u|U) := F_{\alpha}(u) - F_{\alpha}(U) - \nabla F_{\alpha}(U)\nabla A(U)^{-1}(A(u) - A(U))
$$
  
for  $\alpha = 1, ..., d$ .

If we assume  $(H1)$  –  $(H3)$  hold and  $\lim_{|u| \to \infty} \frac{A(u)}{\eta(u)} = 0$  then  $|F_{\alpha}(u|U)| \leq C_D \eta(u|U).$ 

An analogue lemma under more restrictive assumptions

$$
\lim_{|u|\to\infty}\frac{A(u)}{\eta(u)}=\lim_{|u|\to\infty}\frac{F_{\alpha}(u)}{\eta(u)}=0,
$$
\n(2)

was proved in Christoforou & Tzavaras 2017. Note however that this condition is satisfied for polyconvex elastodynamics but is not satisfied e.g. for compressible Euler equations.

#### Theorem

Let  $(\nu, m_A, m_{F_\alpha}, m_\eta)$ ,  $\alpha = 1, ..., d$ , be a dissipative measure-valued solution to [\(1\)](#page-27-0) generated by a sequence of approximate solutions. Let  $U \in W^{1,\infty}(Q)$  be a strong solution to [\(1\)](#page-27-0) with the same initial data  $\eta(u_0) \in L^1(\mathbb{T}^d)$ , thus  $\nu_{0,x} = \delta_{u_0(x)}$ ,  $m_A^0 = m_{F_\alpha}^0 = m_\eta^0 = 0$ . Then  $\nu_{t,x} = \delta_{U(x)}$ ,  $m_A = m_{F_{\alpha}} = m_{\eta} = 0$  and  $u = U$  a.e. in Q.

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## E. Wiedemann.

Weak-strong uniqueness in fluid dynamics, arXiv:1705.04220.

F T. Debiec, P. Gwiazda, K. Lyczek, A. Swierczewska-Gwiazda, ´ Relative entropy method for measure-valued solutions in natural sciences

preprint

## <span id="page-33-0"></span>Thank you for your attention